## Oxygen diffusion and reactions in Hf-based dielectrics

L. V. Goncharova,<sup>a)</sup> M. Dalponte, D. G. Starodub, and T. Gustafsson Department of Physics and Astronomy, and Laboratory for Surface Modification, Rutgers University, 136 Frelinghuysen Road, Piscataway, New Jersey 08854

## E. Garfunkel

Department of Chemistry and Chemical Biology, and Laboratory for Surface Modification, Rutgers University, 610 Taylor Road, Piscataway, New Jersey 08854

P. S. Lysaght, B. Foran, J. Barnett, and G. Bersuker

SEMATECH, 2705 Montopolis Drive, Austin, Texas 78741

(Received 27 March 2006; accepted 25 May 2006; published online 26 July 2006)

Oxygen transport in and reactions with thin hafnium oxide and hafnium silicate films have been investigated using medium energy ion scattering in combination with <sup>18</sup>O<sub>2</sub> isotopic tracing methods. Postgrowth oxidation of Hf-based films in an <sup>18</sup>O<sub>2</sub> atmosphere at 490–950 °C results in O exchange in the film. The exchange rate is faster for pure hafnium oxides than for silicates. The amount of exchanged oxygen increases with temperature and is suppressed by the SiO<sub>2</sub> component. Films annealed prior to oxygen isotope exposure show complex incorporation behavior, which may be attributed to grain boundary defects, and SiO<sub>2</sub> phase segregation. © 2006 American Institute of Physics. [DOI: 10.1063/1.2221522]

Transition metal oxides and silicates with dielectric constant higher than that of SiO<sub>2</sub> are currently being investigated as potential gate dielectric materials in complementary metal-oxide-semiconductor (CMOS) devices.<sup>1</sup> Many fundamental properties of O diffusion in hafnia remain unclear, in particular, the nature of the diffusing oxygen species, the role of oxygen vacancies,<sup>2</sup> interstitial oxygen, and other defects. Previous experiments with ZrO<sub>2</sub> films<sup>3</sup> and ultrafine grained ZrO<sub>2</sub>,<sup>4</sup> and density functional calculations of oxygen incorporation and diffusion energies in monoclinic hafnia<sup> $\circ$ </sup> (HfO<sub>2</sub>) have suggested that oxygen incorporates and diffuses in atomic (ionic, nonmolecular) form. Furthermore, O<sup>2-</sup> becomes a more thermodynamically stable interstitial by accepting two electrons. Calculations<sup>5</sup> show that diffusion via oxygen lattice exchange should be the favored mechanism, however, the barriers for interstitial oxygen diffusion in HfO<sub>2</sub> are small, and defects could be mobile under high temperature processing conditions. In contrast, molecular oxygen incorporation is preferred for the less dense SiO<sub>2</sub> structure, with diffusion proceeding through interstitial sites.<sup>6</sup>

We have studied the diffusion of oxygen in Hf-based dielectric thin films using an isotopic tracing approach. Dielectric films grown with the conventional ( $^{16}O_2$ ) isotope were subsequently annealed in  $^{18}O_2$  (98% isotopically enriched). Medium energy ion scattering (MEIS) was used to measure the  $^{18}O$  and  $^{16}O$  profiles to quantitatively determine the depth distribution of both oxygen species through the dielectric film. Factors in the Hf-based oxides and silicates affecting oxygen exchange<sup>7</sup> such as film composition, phase separation, and crystallinity were examined. It is found that SiO<sub>2</sub> suppresses oxygen diffusion when present on the outer surface of the dielectric. Our data further indicate that HfO<sub>2</sub>–SiO<sub>2</sub> phase separation results in the formation a SiO<sub>2</sub> enriched film close to the top surface of Hf silicate.

2-3 nm Hf oxide and silicate films were deposited on a 1 nm SiO<sub>2</sub>/Si(001) film using atomic layer deposition

<sup>a)</sup>Electronic mail: lgoncha@physics.rutgers.edu

(ALD) at 325 °C with O<sub>3</sub> as an oxidation agent. Reoxidation in <sup>18</sup>O<sub>2</sub> was performed *in situ* in an UHV chamber (~10<sup>-9</sup> Torr) by stabilizing the sample at a temperature in the ~490–950 °C range (measured by an optical pyrometer), followed by <sup>18</sup>O<sub>2</sub> gas introduction at a pressure of 0.01 Torr (30 min). In this letter we focus only on the exchange reaction<sup>7</sup> in the Hf oxide and silicate films.

MEIS was used to determine the depth profile of all elements in the dielectric. We used a H<sup>+</sup> beam with an energy of 130 keV.<sup>8-10</sup> Depth profiles of the elements were obtained by computer simulations of the backscattered ion energy distributions. The depth resolution was  $\sim$ 3 Å in the near surface region and  $\sim$ 8 Å at a depth of 30 Å.

Figure 1(a) shows the part of the backscattered ion spectrum corresponding to the O peak position for the asdeposited HfO<sub>2</sub> film, where the oxygen yield has contributions from both HfO<sub>2</sub> and the SiO<sub>2</sub> interfacial layers. The as-deposited HfO2 films were found to be oxygen-rich  $HfO_{2,10\pm0.05}$ . Strong oxygen exchange is observed in the HfO<sub>2</sub> film at 490 °C (0.01 Torr), as shown in Fig. 1(a). There are no changes in the Si and Hf peaks (not shown), implying that no additional interfacial growth occurs and that atomic O diffusion through the interfacial SiO<sub>2</sub> layer is inefficient under these conditions. Based on a full analysis of the Si, <sup>16</sup>O, <sup>18</sup>O, and Hf peak shapes and energies, we conclude that the interfacial  $SiO_2$  contains only <sup>16</sup>O, as illustrated by the elemental depth distributions in the insert. Excess oxygen atoms, most likely stabilized as interstitials, remain in the hafnia film, as the total number of oxygen atoms remains the same after annealing at 490 °C.

Figure 1(b) shows the MEIS result for  ${}^{18}O_2$  reoxidation when it follows an UHV crystallization annealing at 750 °C for 40 min. [HfO<sub>2</sub> crystallization is expected and was reported at 750 °C (as well as at somewhat higher temperatures for Hf silicates).<sup>11,12</sup>] After  ${}^{18}O_2$  reoxidation of the "crystallized" HfO<sub>2</sub> film, we observe a different oxygen exchange profile, characterized by an  ${}^{18}O$  concentration gradi-

**89**, 044108-1

<sup>© 2006</sup> American Institute of Physics

Downloaded 16 Sep 2008 to 129.100.41.190. Redistribution subject to AIP license or copyright; see http://apl.aip.org/apl/copyright.jsp



FIG. 1. The <sup>16</sup>O and <sup>18</sup>O energy range of MEIS spectra for HfO<sub>2</sub> films on Si (100) before (open circles) and after oxidation in <sup>18</sup>O<sub>2</sub> (solid symbols) at 490 °C. Deconvolutions of the peak into HfO<sub>2</sub> and SiO<sub>2</sub> contributions are shown by the thin lines. The spectra were obtained with 130.3 keV protons in a channeling and blocking configuration. The HfO<sub>2</sub> film is amorphous in (a) and crystalline in (b). Inserts show best fit elemental depth distribution profiles for the samples after <sup>18</sup>O oxidation.

ent decreasing from the top surface towards the interface with the substrate [insert, Fig. 1(b)].

Recent high resolution transmission electron microscopy (HRTEM) studies of  $HfO_2$  films<sup>2</sup> indicated that there is O depletion in the grain boundary region after rapid thermal annealing. In other works on the analogous  $ZrO_2$  case, faster diffusion was found to occur via the grain boundaries of nanocrystalline films, compared to the "bulk" of nanocrystallites, and was attributed to higher vacancy concentration in the oxygen sublattice in the grain boundary region.<sup>4</sup> Consolidation of nanocrystalline grains in as-deposited HfO<sub>2</sub> films into larger-size crystallites after vacuum anneals results in fewer grain boundaries. Therefore, the contribution of the bulk of the crystallites with slower diffusion becomes more prominent. The observed <sup>18</sup>O distribution can be considered characteristic of bulk HfO<sub>2</sub> crystallite (major fraction, slow) and grain boundary diffusion (minor, fast).

Several factors may also contribute to different exchanged oxygen distributions in the as-deposited and crystallized films. Presuming that the transported species are individual oxygen atoms (ions, not molecules), the availability of atomic oxygen at the surface is one of the factors affecting Downloaded 16 Sep 2008 to 129 100 41 190 Redistribution suble

TABLE I. Areal densities of different oxygen isotopes before and after annealing in an  ${}^{18}O_2$  atmosphere ( $p_{18_{O_2}}$ =10<sup>-2</sup> Torr) at 490 °C for 30 min.

Sample	Total oxygen $(\times 10^{15} \text{ at./cm}^2)$	$^{16}O$ loss ( $^{18}O$ gain) (×10 $^{15}$ at./cm $^{2}$ )	Exchange fraction
HfO <sub>2</sub>	15.6	7.8(7.8)	0.50
(HfO <sub>2</sub> ) <sub>2</sub> SiO <sub>2</sub>	13.3	2.6(2.6)	0.20
HfO <sub>2</sub> SiO <sub>2</sub>	14	< 0.5(0.5)	< 0.04
HfO <sub>2</sub> (crystalline)	14.7	5.6(5.8)	0.40
(HfO <sub>2</sub> ) <sub>2</sub> SiO <sub>2</sub> (crystalline)	13.2	0.9-1.5(1.1)	0.07-0.11
$SiO_2/HfO_2$	19.5	1.0(1.0)	0.07

the extent of the exchange in the oxide.  $O_2$  is expected to be adsorbed molecularly on perfect surfaces of HfO<sub>2</sub> and to dissociate primarily at O-vacancy defect sites.<sup>13</sup> The amount of available atomic O depends on the rate of  $O_2$  dissociation at the surface and is therefore related to the number of oxygen vacancies at the top surface. Additionally, as relaxation of atoms along the diffusion path is important for a lattice exchange mechanism, as-deposited disordered HfO<sub>2</sub> films might be expected to display a lower diffusion barrier than crystalline films.

In order to quantify and discuss the amount of <sup>18</sup>O incorporated into the high- $\kappa$  layer, we calculate the <sup>18</sup>O exchange fraction as the ratio of <sup>18</sup>O to the total oxygen (<sup>16</sup>O + <sup>18</sup>O) MEIS areal density in an  $n(\text{HfO}_2)m(\text{SiO}_2)$  film (excluding any interfacial SiO<sub>2</sub>). Representative exchange fractions for selected as-deposited and crystalline Hf oxide and silicate films are listed in Table I. This comparison shows that when SiO<sub>2</sub> is introduced into HfO<sub>2</sub> films during the growth, the exchange process slows dramatically and is almost completely suppressed for 50% SiO<sub>2</sub> content (under our reoxidation conditions: 490 °C, 10<sup>-2</sup> Torr, and 30 min).

Upon vacuum annealing to  $\geq 850$  °C some changes were observed in the Hf and Si peak shapes (Fig. 2) and therefore in their depth distributions [in the (HfO<sub>2</sub>)<sub>2</sub>SiO<sub>2</sub> film]. The Hf peak increases in height and decreases in width, with the Hf areal density remains constant at (3.60±0.15)×10<sup>15</sup> Hf atoms/cm<sup>2</sup>. While the Si distribution is uniform in as-deposited (HfO<sub>2</sub>)<sub>2</sub>SiO<sub>2</sub>, after annealing the



FIG. 2. The Hf and Si MEIS peaks for a  $(HfO_2)_2SiO_2$  film before (solid line) and after annealing in vacuum for 30 min (850 °C, 10<sup>-9</sup> Torr) (circles). Simulation profiles for Hf, Si, and O distributions for the annealed film are shown in the insets.

Downloaded 16 Sep 2008 to 129.100.41.190. Redistribution subject to AIP license or copyright; see http://apl.aip.org/apl/copyright.jsp



FIG. 3. The <sup>18</sup>O exchange fraction in a 29 Å thick  $(HfO_2)_2SiO_2$  film as a function of oxidation temperature (30 min,  $p_{18}_{O_2} = 10^{-2}$  Torr) for asdeposited (open circles) and crystallized or phase segregated (dark circles) films.

Si yield becomes larger closer to the top surface. This change can be a signature of phase segregation of the silicate layer with SiO<sub>2</sub> enrichment in the top surface layer. The corresponding depth profile (insert) is illustrative only, because strong lateral inhomogeneities are expected in this case. After reoxidation at 490 °C, only the O exchange reaction is observed, with no observable net increase in oxygen. The oxygen exchange fraction for recrystallized silicate is in the 0.07-0.11 range<sup>14</sup> and is lower than for the as-deposited film. Since a crystallization annealing should lead to phase segregation in the Hf silicates, the lower O exchange fraction results from a lower surface area of HfO<sub>2</sub> exposed to oxygen, or a suppression of the diffusion through the silica enriched grain boundary regions.

Figure 3 shows the increase of the <sup>18</sup>O exchange fraction in a (HfO<sub>2</sub>)<sub>2</sub>SiO<sub>2</sub> film as a function of reoxidation temperature in the 490–950 °C range. This exchange fraction increase is driven by several factors. First, the <sup>18</sup>O<sub>2</sub> dissociation rate at the surface goes up as the temperature increases. Second, more oxygen species will diffuse towards the interface and will be trapped there due to chemical reactions with the Si substrate (if there is no barrier layer at the interface). Only exchange is observed at 490 °C; at higher temperatures ( $\geq$ 610 °C), interfacial SiO<sub>2</sub> growth is apparent, as others have observed under similar conditions.<sup>3</sup>

In order to understand the suppression of the O exchange in Hf silicates after crystallization annealing, and how it affects O<sub>2</sub> diffusion and dissociation, we deposited an ~2 ML thick SiO<sub>2</sub> layer on top of the HfO<sub>2</sub> film. The ~2 ML SiO<sub>2</sub> on top of the HfO<sub>2</sub> film was achieved by two SiO<sub>2</sub> ALD deposition cycles at 410 °C. We find that the exchange rate at 490 °C is reduced to almost zero (last line, Table I) indicating that the oxygen diffusion is stopped by the thin SiO<sub>2</sub> layer. This is consistent with our results in Fig. 1, which shows no exchange in the interfacial SiO<sub>2</sub> film while atomic O is in abundance in the adjacent hafnia film. Alternatively, the top SiO<sub>2</sub> layer is expected to suppress O<sub>2</sub> dissociation by eliminating Hf bonded defect sites on the top surface thus reducing supply of atomic O.

These data suggest that the suppression of O exchange in Hf silicates discussed above may be caused by the segregation of  $SiO_2$  to the top surface that accompanies the phase separation which is known to occur in Hf silicates under the high temperature annealing.<sup>12,15</sup> Indeed, as both Si–O and Hf–O bonds are present on the surface of Hf silicates, the

<sup>18</sup>O<sub>2</sub> dissociation rate is expected to be lower for silicates with higher SiO<sub>2</sub> content. Nevertheless, within a simple assumption that the Hf–O/Si–O bond ratio at the surface is the same as in the  $n(HfO_2)m(SiO_2)$  film bulk, the exchange rate should drop by less than a factor of 2 by going from HfO<sub>2</sub> to HfO<sub>2</sub>SiO<sub>2</sub>. However, a much more dramatic exchange rate decrease points to a nonuniform Hf:Si surface distribution with a higher SiO<sub>2</sub> content (SiO<sub>2</sub> segregation) close to the top surface, probably due to its lower surface energy.<sup>16,17</sup> This surface layer enrichment with SiO<sub>2</sub> can significantly reduce and even block O<sub>2</sub> diffusion at low temperatures. A second effect is the suppression of oxygen lattice exchange in the Hf silicate due to the presence covalent Si–O bonding, as compared to the more ionic hafnium oxide.

In summary, we have studied oxygen exchange in Hfbased dielectric films by high-resolution ion profiling with isotopic tracing. Our results indicate that atomic oxygen diffusion via oxygen lattice exchange is the predominant diffusion mechanism in hafnia, consistent with theoretical calculations.<sup>5</sup> We further show that the exchange rate scales with temperature. Finally, we show that the addition (and surface segregation) of SiO<sub>2</sub> to hafnium oxide and Hf silicate suppresses O incorporation in the dielectric.

The authors thank the SRC/Sematech FEP-TC, the NSF, and CAPES (Brazil) for financial support.

- <sup>1</sup>J. J. Peterson, C. D. Young, J. Barnett, S. Gopalan, J. Gutt, C.-H. Lee, H.-J. Li, T.-H. Hou, Y. Kim, C. Lim, N. Chaudbary, N. Moumen, B.-H. Lee, G. Bersuker, G. A. Brown, P. M. Zeitzoff, M. I. Gardner, R. W. Murto, and H. R. Huff, Electrochem. Solid-State Lett. **7**, G164 (2004).
- <sup>2</sup>H. S. Baik, M. Kim, G.-S. Park, S. A. Song, M. Varela, A. Franceschetti, S. T. Pantelides, and S. J. Pennycook, Appl. Phys. Lett. **85**, 672 (2004).
  <sup>3</sup>B. W. Busch, W. H. Schulte, E. Garfunkel, T. Gustafsson, W. Qi, R. Nieh, and J. Lee, Phys. Rev. B **62**, R13290 (2002).
- <sup>4</sup>U. Brossmann, R. Wurschum, U. Sodervall, and H. Schaefer, J. Appl. Phys. **85**, 7646 (1999).
- <sup>5</sup>A. S. Foster, A. L. Shluger, and R. M. Nieminen, Phys. Rev. Lett. **89**, 225901 (2002).
- <sup>6</sup>S. Mukhopadhyay, P. Sushko, A. M. Stoneham, and A. L. Shluger, Phys. Rev. B **71**, 235204 (2005).
- <sup>4</sup>We define oxygen exchange as substitution of the oxygen atoms within the framework of the thin film by oxygen coming from the gas phase. The number of atoms gained from the gas phase equals to the number of oxygen atoms lost, so the total oxygen content of the dielectric remains constant during exchange. This exchange process has to be differentiated from interfacial incorporation and growth, where oxygen reacts with the film or the substrate, increasing the total oxygen content.
- <sup>8</sup>W. H. Schulte, B. W. Busch, E. Garfunkel, T. Gustafsson, G. Schiwietz, and P. L. Grande, Nucl. Instrum. Methods Phys. Res. B **183**, 16 (2001). <sup>9</sup>R. M. Tromp, M. Copel, M. C. Reuter, M. Horn von Hoegen, J. Speidell,
- and R. Koudijs, Rev. Sci. Instrum. **62**, 2679 (1991).
- <sup>10</sup>J. F. Ziegler and J. P. Biersack, sRIM-the Stopping and Range of Ions in Matter, Version 2003.20, 2003.
- <sup>11</sup>S. V. Ushakov, A. Navrotsky, Y. Yang, S. Stemmer, K. Kukli, M. Ritala, M. A. Leskela, P. Fejes, A. A. Demkov, C. Wang, B. Y. Nguyen, D. Triyoso, and P. Tobin, Phys. Status Solidi B **241**, 2268 (2004).
- <sup>12</sup>S. Stemmer, Y. Li, B. Foran, P. S. Lysaght, S. Streiffer, P. Fuoss, and S. Seifert, Appl. Phys. Lett. 83, 3141 (2003).
- <sup>13</sup>V. E. Henrich and P. A. Cox, *The Surface Science of Metal Oxides* (Cambridge University Press, Cambridge, 1994).
- <sup>14</sup>The fluctuations in the exchange fraction for crystalline films may come from sample-to-sample variations in the crystallite size distribution or from surface SiO<sub>2</sub> segregation.
- <sup>15</sup>H. Kim and P. C. McIntyre, J. Appl. Phys. **92**, 5094 (2002).
- <sup>16</sup>E. Chagarov, A. A. Demkov, and J. B. Adams, Phys. Rev. B **71**, 075417 (2005).
- <sup>17</sup>A. A. Demkov (private communication).